Simulating Quantum Field Theories on a Quantum Computer

Stephen Jordan







Can quantum computers simulate all physical processes efficiently?

Universality Conjecture:

Quantum circuits can simulate all physical dynamics in $poly(E, V, t, 1/\epsilon)$ time.

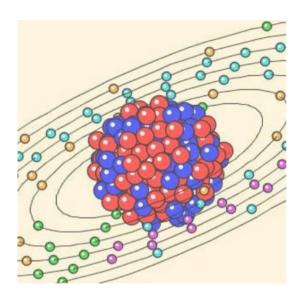
Status:

Non-relativistic QM	Yes: Now being optimized	
Quantum Field Theories	Probably: In progress	
Quantum Gravity/Strings	Nobody knows	

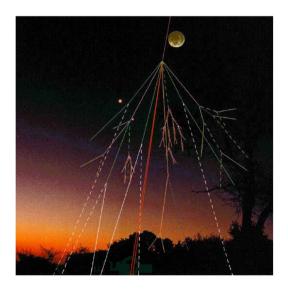
Quantum Field Theory

- Much is known about using quantum computers to simulate quantum systems.
- Why might quantum field theory be different?
 - Field has infinitely many degrees of freedom
 - Relativistic
 - Particle number not conserved
 - Formalism looks different.

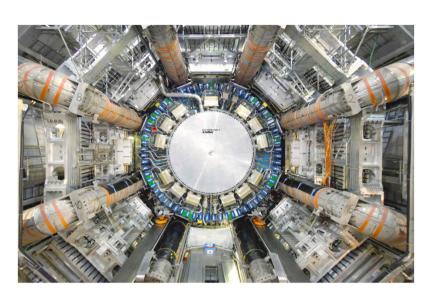
When do we need QFT?



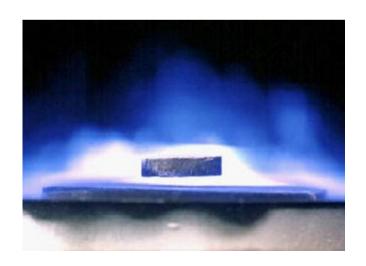
Nuclear Physics



Cosmic Rays



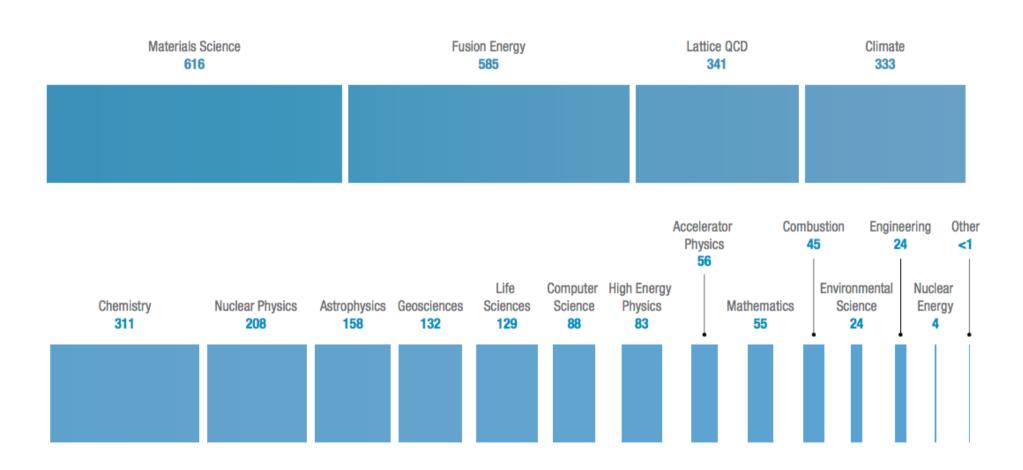
Accelerator Experiments



Coarse-grained many-body systems

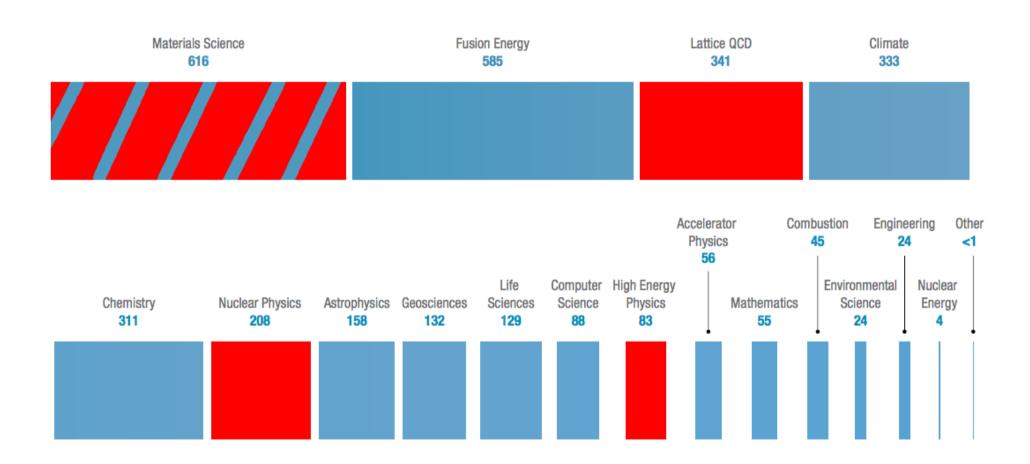
2015 NERSC Usage By Discipline

(MPP HOURS IN MILLIONS)



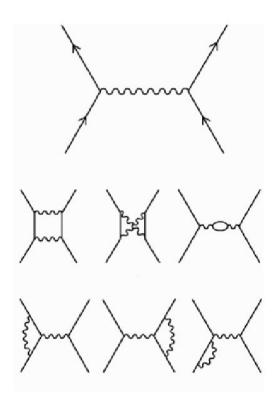
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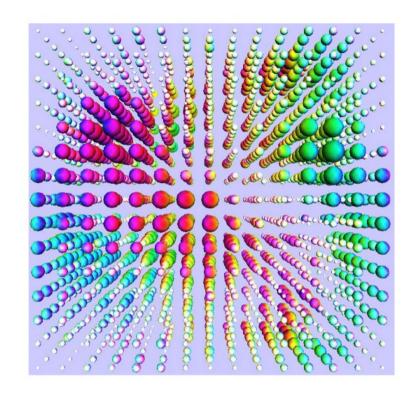
Classical Algorithms

Feynman diagrams



Break down at strong coupling or high precision

Lattice methods



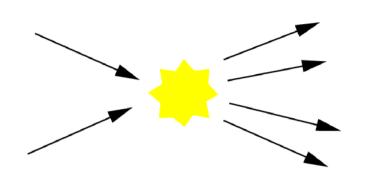
Cannot calculate scattering amplitudes

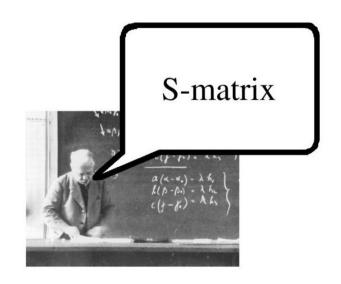
There's room for exponential speedup by quantum computing.

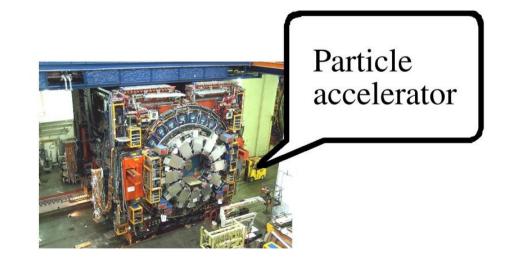
A QFT Computational Problem

Input: a list of momenta of incoming particles.

Output: a list of momenta of outgoing particles.







Results So Far

Efficient quantum simulation algorithms:

	Bosonic	Fermionic	
Massive	Jordan, Lee, Preskill Science, 336:1130 (2012)	Jordan, Lee, Preskill <i>ArXiv:1404.7115</i> (2014)	
Massless	?	?	

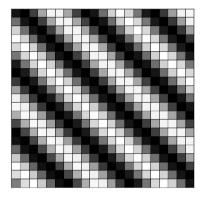
 BQP-hardness: classical computers cannot perform certain QFT simulations efficiently

[S. Jordan, H. Krovi, K. Lee, K. Preskill, 2017]

Better Speed and Broken Symmetries
 [A. Moosavian and S. Jordan, 2017]

Representing Quantum Fields

A field is a list of values, one for each location in space.



A quantum field is a superposition over classical fields.

$$\frac{1}{\sqrt{2}}$$
 $\left| \frac{i}{\sqrt{2}} \right|$

A superposition over bit strings is a state of a quantum computer.

ϕ^4 -theory

Lagrangian density

$$\mathcal{L} = \frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi - m^2 \phi^2 - \frac{\lambda}{4!} \phi^4$$

For quantum simulation we prefer Hamiltonian formulation (equivalent)

$$H = \int d^dx \left[\pi^2 + (\nabla \phi)^2 + m^2 \phi^2 + \lambda \phi^4 \right]$$

$$[\phi(x), \pi(y)] = i\delta^{(d)}(x - y)$$

Our Algorithms

1) Choose a lattice discretization.

Bound discretization error (by renormalization group).

2) Prepare physically realistic initial state.

Is the most time-consuming step.

This depends strongly on which QFT simulated.

3) Implement time-evolution by a quantum circuit.

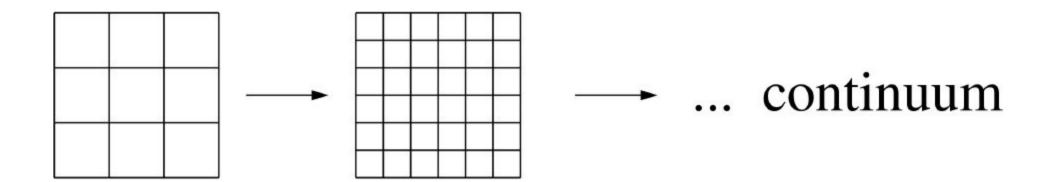
Use Trotter formulae.

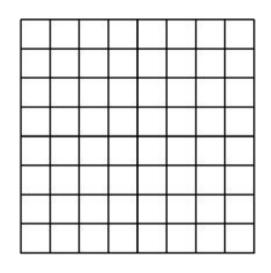
4) Perform measurements on final state.

Complicated by vacuum entanglement.

Lattice Cutoff

Continuum QFT = limit of a sequence of theories on successively finer lattices.

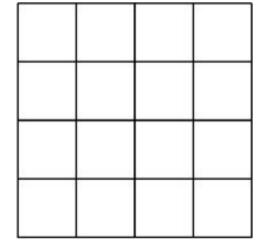




Mass: m

Interaction strength: λ

Coarse grain

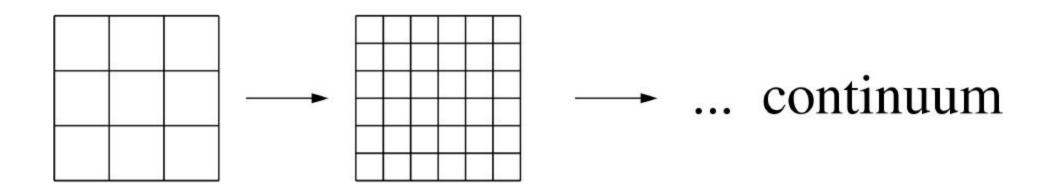


Mass: m'

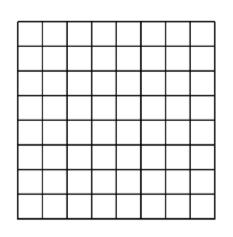
Interaction strength: λ'

Lattice Cutoff

Continuum QFT = limit of a sequence of theories on successively finer lattices.



m and λ are functions of lattice spacing!



$$H = \frac{1}{2} \sum_{x \in \Omega} a^d \left[\pi^2 + (\nabla \phi)^2 + m^2 \phi^2 + \lambda \phi^4 \right]$$

Coarse grain

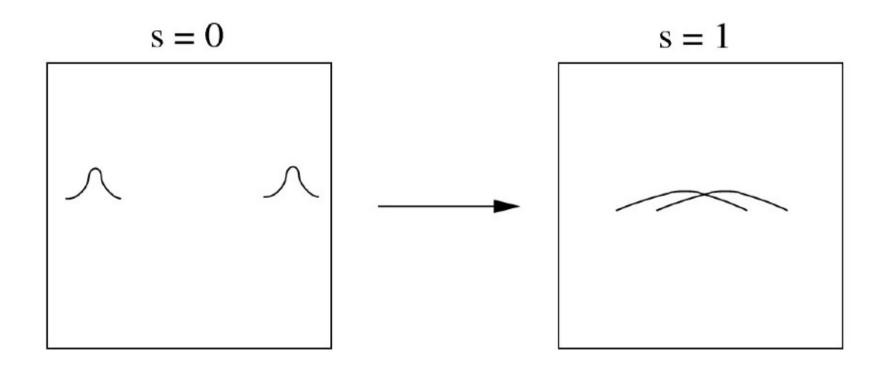
$$H_{\text{eff}} = \frac{1}{2} \sum_{x \in \Omega'} (2a)^d \left[\pi^2 + (\nabla' \phi)^2 + m_{\text{eff}}^2 \phi^2 + \lambda_{\text{eff}} \phi^4 + g \phi^6 + \dots \right]$$

Simulation converges as a^2

Adiabatic State Preparation

$$H(s) = H_{\text{free}} + sH_{\text{interaction}}$$

Prepare wavepackets in free theory, then adiabatically turn on interaction. Problem:

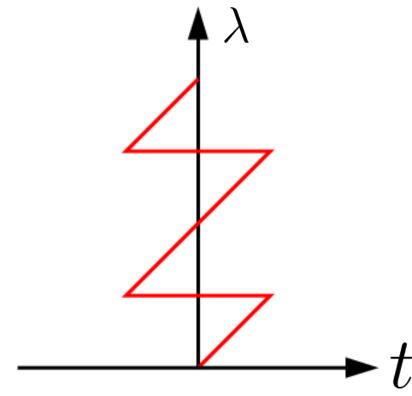


Adiabatic State Preparation

Solution: intersperse backward time evolutions with time-independent Hamiltonians.

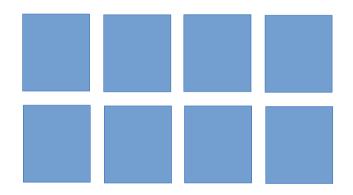
This winds back dynamical phase on each eigenstate without undoing adiabatic change of

basis.



Simulating Detectors

Measure energy in localized regions:



 Need smooth envelope function to avoid excessive vacuum noise!

$$H_f = \sum_{\mathbf{x}} f(\mathbf{x}) \mathcal{H}(\mathbf{x})$$

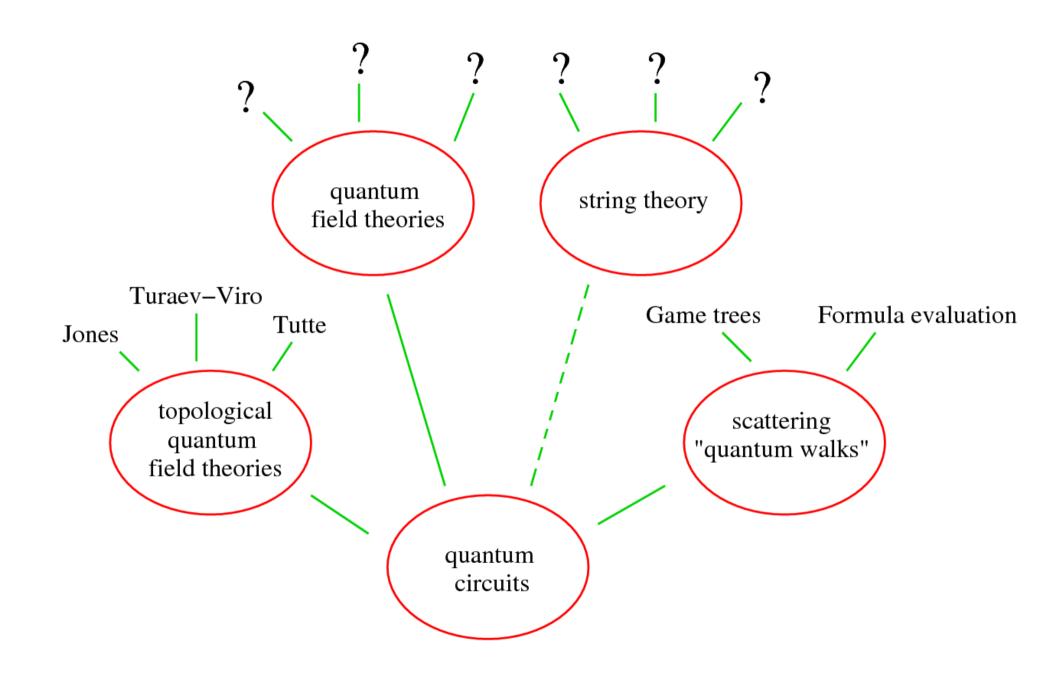
Runtimes

Weak Coupling:

d=1	$(1/\epsilon)^{1.5}$
d=2	$(1/\epsilon)^{2.376}$
d=3	$(1/\epsilon)^{5.5}$

Strong Coupling:

	$\lambda_c - \lambda_0$	p	$n_{ m out}$
d = 1	$\left(\frac{1}{\lambda_c - \lambda_0}\right)^9$	p^4	$n_{ m out}^5$
d=2	$\left(\frac{1}{\lambda_c - \lambda_0}\right)^{6.3}$	p^6	$n_{ m out}^{7.128}$



Fermions:

- Fermion doubling problem
- Free vacuum different from Bosonic case

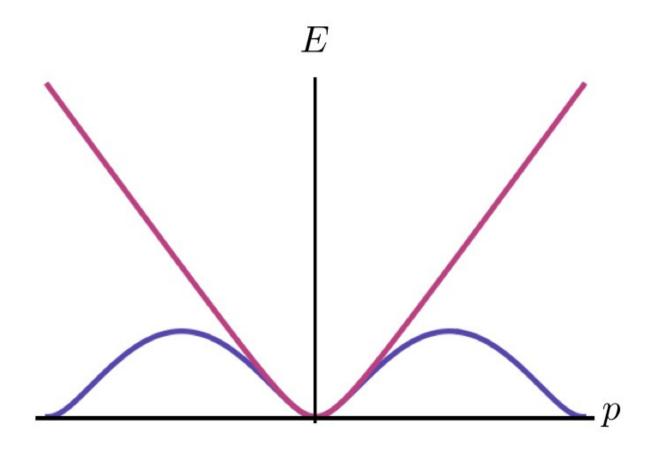
Gross-Neveu:

$$H = \int dx \left[\sum_{j=1}^{N} \bar{\psi}_j \left(m_0 - i\gamma^1 \frac{d}{dx} \right) \psi_j + \frac{g^2}{2} \left(\sum_{j=1}^{N} \bar{\psi}_j \psi_j \right)^2 \right]$$

Fermion Doubling Problem

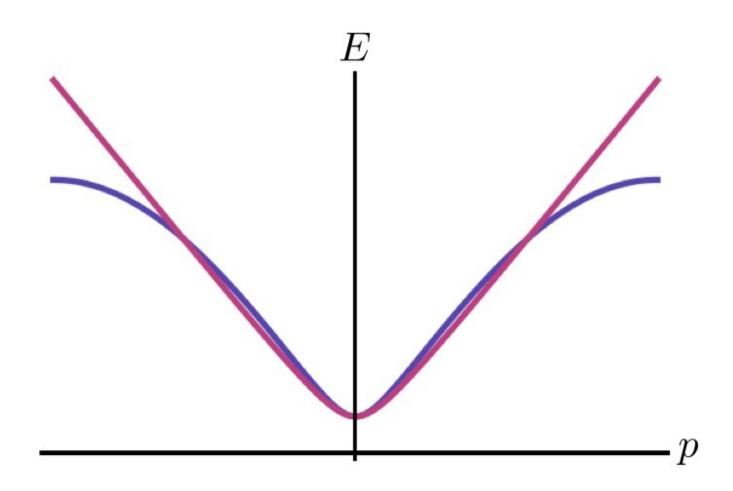
$$\frac{d\psi}{dx} \to \frac{\psi(x+a) - \psi(x-a)}{2a}$$

$$\sqrt{p^2 + m^2} \to \sqrt{\sin^2 p + m^2}$$



Wilson Term

$$H \to H - \frac{r}{2a} \sum_{x} \bar{\psi} \left(\psi(x+a) - 2\psi(x) + \psi(x-a) \right)$$



Improved State Prep: Bosons

- In some cases (e.g. weakly coupled d=2),
 preparing the free vacuum is the rate limiting step.
- We can do this much faster using Bogoliubov transformation that looks like a Fast Fourier Transform.

[Somma, Jordan, unpublished]

Essentially same idea as 2nd quantized FFT from:

[Babbush, Wiebe, McClean, McLain, Neven, Chan, 2017]

Improved State Prep: Fermions

- Two problems with adiabatic state preparation:
 - Cannot reach symmetry-broken phase
 - Runtime bound not practical: $O(\epsilon^{-8})$
- A solution for both:
 - First, prepare the vacuum from MPS
 - Then, resonantly excite single-particle wavepackets
 - Tighter analysis: CFT entropy and Floquet theory:

$$O(\epsilon^{-3.23})$$

[A. Moosavian, S. Jordan, 2017]

Tensor Network Ansatzes

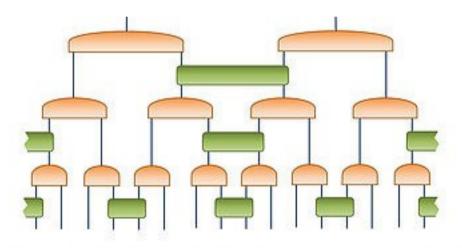
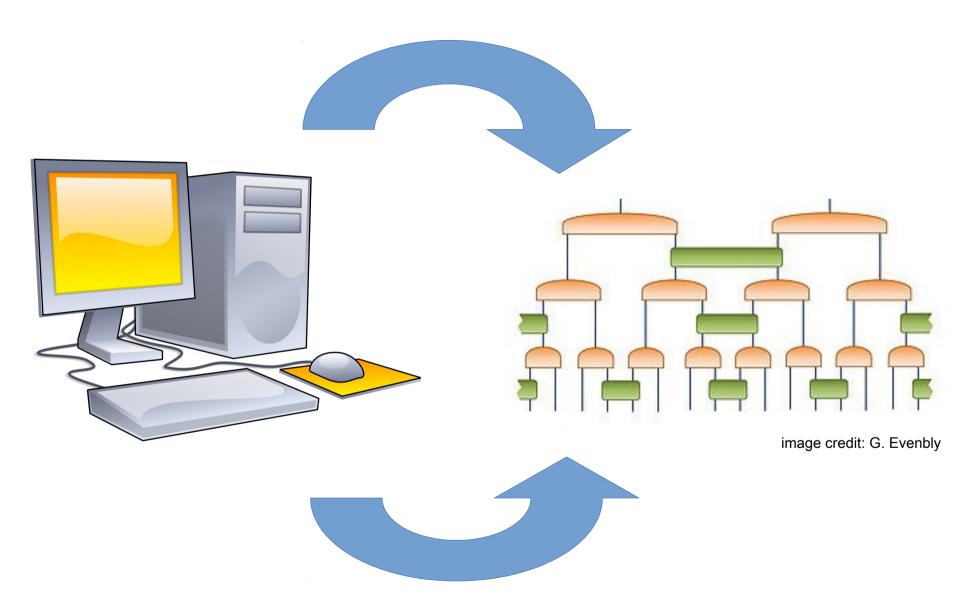


image credit: G. Evenbly

Tensor Network Ansatzes



Near-Term Prospects?

quantum supremacy

science applications

commercial applications

- Simulating conformal field theories using MERA-based variational eigensolvers
- Simulating commuting Hamiltonians
- Simulating high-connectivity systems, e.g. spin glasses or SYK model

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Thanks!